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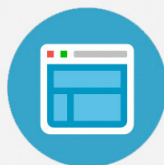
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# Silicon photodetectors integrated with vertical silicon nitride waveguides as image sensor pixels: Fabrication and characterization

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The current trend toward image sensors with ever-increasing pixel counts is prompting continual reductions in pixel area, leading to significant cross-talk and efficiency challenges. The realization of image sensor pixels containing waveguides presents a means for addressing these issues. The fabrication of such pixels is however not straightforward. Conventional waveguides employed in integrated optics are horizontal, but waveguides needed for the proposed sensor must be vertical and integrated with photodetectors. Here, the authors describe a fabrication process for vertical silicon nitride waveguides integrated with silicon photodetectors. The authors describe the etching, deposition, and planarization techniques that enable the formation of silicon nitride waveguides embedded in silicon dioxide. They also describe a fabrication process for silicon photodetectors, including a means for ensuring that their photosensitive areas have sizes consistent with those of photodetectors employed in conventional image sensors. In addition, the authors perform optical and electrical characterization of the fabricated devices. The results demonstrate the ability of the fabricated waveguides to guide light onto the photodetectors with high efficiency. © 2014 American Vacuum Society. [<http://dx.doi.org/10.1116/1.4868627>]

## I. INTRODUCTION

In recent years, there has been enormous growth in the prevalence of digital imaging systems in consumer applications. The image sensors used in these systems are being increasingly dominated by complementary metal oxide semiconductor (CMOS) technology. For many applications, there has been a market demand for increased spatial resolution. This has been largely achieved by decreasing the pixel size, which has reduced from more than 10  $\mu\text{m}$  to less than 2  $\mu\text{m}$  over approximately one decade.<sup>1</sup> It has been predicted that optical efficiency will be reduced and interpixel cross-talk

increased as conventional image sensor pixels are scaled to submicron dimensions.<sup>1</sup> As we discuss further below, vertical waveguides, or “light pipes,” present a means for overcoming these fundamental challenges. Fesenmeier *et al.* simulated light pipe structures.<sup>2</sup> Hsu *et al.*<sup>3</sup> and Gambino *et al.*<sup>4</sup> demonstrated light pipe devices experimentally, with pipe heights ranging from 2.2 to 4  $\mu\text{m}$ , by stacking multiple etched sections. We previously demonstrated the ability of silicon nitride light pipes to improve optical efficiency.<sup>5</sup> Here, we describe the fabrication processes that enables the formation of silicon nitride light pipes integrated with photodetectors. These are described in a level of detail not present in Ref. 5. We demonstrate the deposition of silicon nitride with good optical properties and the etching process to form pillars. We

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also present a process for encasing the pillars in silicon dioxide by deposition and chemical mechanical polishing (CMP). The integration of the light pipes formed in this way with silicon photodetectors is discussed. We demonstrate a means for ensuring that the photosensitive regions of these photodetectors have sizes comparable to those of traditional image sensor pixels. As discussed below, this is critical in order for optical performance enhancement enabled by our light pipes to be demonstrated. The current–voltage characteristics of the completed devices under dark and illuminated conditions are presented. Finally, we present results confirming that the light pipes guide light onto the photodetectors with high efficiency.

The ability of light pipes to improve optical performance in scaled image sensor pixels arises from both geometric optics and diffraction effects. We begin by considering geometric optics effects. Light incident on a conventional image sensor pixel is focused by a microlens through the dielectric stack onto a photodetector. Reducing the pixel size while maintaining the microlens-to-photodetector distance leads to increased interpixel cross talk and lower optical efficiency, especially for light incident at other than normal incidence. This can be understood from Fig. 1. A conventional image sensor pixel is shown as Fig. 1(a). Light at normal incidence is focused to the photodetector center. A consequence of the finite size of the photodetector (diameter  $R$ ), however, is that only light incident on the microlens up to a maximum angle of  $\phi_1$  can be collected onto the photodetector. This angle is given approximately by  $\phi_1 \approx nR/(2F_1)$ . Here,  $n$  is the microlens refractive index (assumed to be equal to that of the dielectric stack) and  $F_1$  is its focal length. As pixel dimensions (and therefore  $R$ ) shrink, and the microlens-to-photodetector distance (and hence  $F$ ) is maintained or increased, the maximum acceptance angle ( $\phi_1$ ) decreases, leading to efficiency and interpixel cross-talk issues. In Fig. 1(b), it can be seen that the light pipe increases the acceptance angle (to  $\phi_2 \approx nR/(2F_2)$ ), due to the fact that the microlens has a shorter focal length ( $F_2$ ). These geometric optics considerations therefore predict that light pipes enhance the optical performance. They also emphasize if a device is built to demonstrate the improved performance possible with light

pipes, then the photodetector will need to have a photosensitive region (diameter  $R$ ) that is limited in extent (and consistent with the dimensions of photodetectors in actual image sensors) for the enhancement to be confirmed. It has also been noted that diffraction effects become very pronounced for submicron pixels,<sup>1</sup> a consequence of the fact that the microlens has an extent of only a few wavelengths, which prohibits efficient focusing of the incident light on the photodetector.<sup>1</sup> Here, light pipes present a means for improving optical performance as the light only needs to be focused onto the light pipe's entrance, rather than through the full dielectric stack onto the photodetector. The use of the light pipe is therefore akin to reducing the stack height. Simulations have predicted that stack height reduction should enable submicron pixels to achieve good optical performance.<sup>1</sup>

The organization of this paper is as follows. We describe the fabrication of the silicon nitride light pipes and silicon photodetectors in Secs. II A and II B, respectively. For the latter, we include a discussion of the fabrication steps that enable the successful integration of light pipes and photodetectors. In Sec. III, optical and electrical measurements are presented to demonstrate that the device achieves the desired efficiency improvement. Conclusions are presented in Sec. IV.

## II. FABRICATION RESULTS

### A. Fabrication of silicon nitride light pipes

One approach to the realization of silicon nitride light pipes embedded in silicon dioxide involves the following steps. Silicon dioxide is first deposited, for example by plasma-enhanced chemical vapor deposition (PECVD). Holes are then produced by standard optical lithography and reactive ion etching. The holes are then filled by the deposition of silicon nitride by PECVD. We have previously demonstrated this approach<sup>6</sup> with challenges. While silicon nitride generally deposits much more conformally by PECVD than by low pressure chemical vapor deposition, it is however not close to being ideal. Deposition rates are higher at the trench corners than those along the trench sides. As a result, the top openings of the holes close, preventing further filling, leading to void formation (Fig. 2). We therefore develop an alternative method

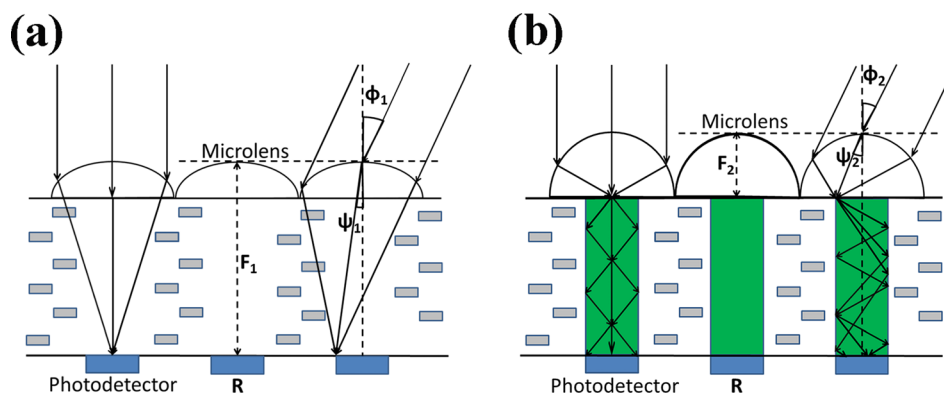


FIG. 1. (Color online) (a) On the left, a typical image sensor pixel where microlens focuses incident light on the center of the photodetector. On the right,  $\Phi_1$  is the maximum incident angle of light that photodetector can detect (b) On the left, an image sensor pixel with light-pipe (waveguide situated between color filter and photodetector) where light is focused by microlens on top of the light-pipe and guided through the photodetector. On the right,  $\Phi_2$  is the maximum incident angle of light that can be focused by microlens on top of the light-pipe and detected by the photodetector at the bottom.

in which the silicon nitride is etched to form pillars, followed by being encased in silicon dioxide. The first step of this process involves the deposition of silicon nitride by PECVD. For the proposed application, it is of course important that the silicon nitride exhibits good optical transmission. The PECVD-grown  $\text{SiN}_x$  film under study has acceptable optical transmission in 400–900 nm spectral range. Figure 3(a) shows the transmission spectra of 9  $\mu\text{m}$  thick films of  $\text{SiN}_x$  and  $\text{SiO}_2$ . It can be seen that  $\text{SiN}_x$  film deposited with a flow rate of 35 sccm  $\text{SiH}_4$  has far higher optical transmission than that deposited at 40 sccm, which is consistent with previous studies.<sup>7,8</sup> In the light pipe devices described in the remainder of this paper, the deposition of  $\text{SiN}_x$  follows the process parameters with the flow rates of  $\text{SiH}_4$ ,  $\text{NH}_3$ , and  $\text{N}_2$  being set to 35, 40, and 1960 sccm, respectively. The radiofrequency (RF) powers at high and low frequencies are 21 W and 50 W, respectively. The platen and lid temperatures are 300°C and 250°C, respectively. The chamber pressure is 900 mTorr. This results in an  $\text{SiN}_x$  deposition rate of 11.6 nm/min. A deposition rate of 18.75 nm/min results when the  $\text{SiH}_4$  flow rate is increased to 40 sccm, while keeping all other parameters the same.

We now describe the fabrication of silicon nitride pillars. We first deposit  $\text{SiN}_x$  to a thickness of 8  $\mu\text{m}$  using the method described above. We then perform optical lithography, thermally evaporate a layer of Al (50 nm) and perform the lift off process. This yields an array of Al disks (5  $\mu\text{m}$  diameter) that serve as etch masks. We consider the etching step we perform next to be the most challenging part of the fabrication process. Inductively coupled plasma reactive ion etching [ICP-RIE, surface technology systems (STS)] is used with an  $\text{SF}_6/\text{C}_4\text{F}_8$  gas mixture. The  $\text{SF}_6$  gas is responsible for the etching, and acts in an isotropic fashion. The  $\text{C}_4\text{F}_8$  gas is responsible for the formation of a polymer layer, and therefore promotes anisotropic etching. We find that it also enhances the smoothness of the sidewalls and of the bottom surface but decreases the etch rate. The etch rate depends on the diameter of the pillars, the distance between the pillars, the ratio of the flow rates of the gases and the pressure in the chamber, the coil, and RF powers. In this work, we experiment with varying the ratio between the flow

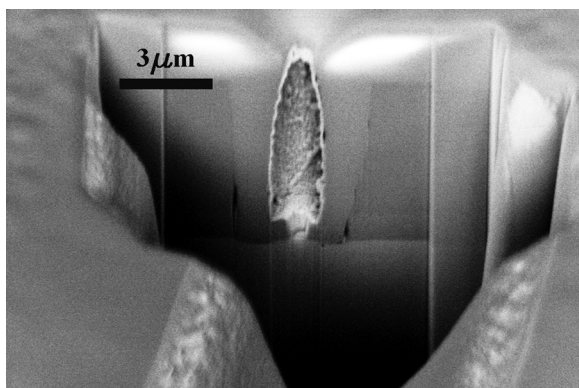


Fig. 2. Scanning electron micrograph of structure resulting when  $\text{SiN}_x$  is deposited into a cylindrical hole opened in a thick  $\text{SiO}_2$  layer. It can be seen that the nonconformal deposition of  $\text{SiN}_x$  results in the formation of a key-hole, i.e., incomplete filling.

rates of the etching gases  $\text{C}_4\text{F}_8$  and  $\text{SF}_6$ , with the aim of achieving silicon nitride pillars with vertical side walls. The  $\text{C}_4\text{F}_8$  flow rate is maintained at 150 sccm for  $\text{C}_4\text{F}_8/\text{SF}_6$  gas flow ratios greater than 1.0, while the  $\text{SF}_6$  flow rate is maintained at 130 sccm for  $\text{C}_4\text{F}_8/\text{SF}_6$  ratios below 1.0. We find that vertical light pipes, with sidewall angles close to 90°, are achieved at relatively high  $\text{C}_4\text{F}_8$  to  $\text{SF}_6$  flow rate ratios [Fig. 3(b)]. The etch rate is reduced at such ratios, but sufficiently high for the light pipes to be produced in a reasonable time. We note however that polymer forms during the etch process, which may lead to micromasking and increased surface roughness. We address this problem by two methods that use  $\text{O}_2$  gas. In one approach, a small amount of  $\text{O}_2$  (4 sccm flow rate) is added to the  $\text{C}_4\text{F}_8/\text{SF}_6$  etch process. This helps remove polymer that forms on the sidewalls. The second method comprises the use of an  $\text{O}_2$  plasma for 30 min after the formation of the light pipes. The use of  $\text{O}_2$  to mitigate the polymer formation is described further below.

We find that the silicon nitride pillars are etched with undercutting of the Al disk masks. The desired final pillar diameter can be achieved by increasing the diameter of the etch mask for the Al disk. Figure 4(a) shows a silicon nitride pillar produced by etching for 2 hours ( $\text{C}_4\text{F}_8$  flow rate: 150 sccm,  $\text{SF}_6$  flow rate: 75 sccm, coil power: 1200 W, platen

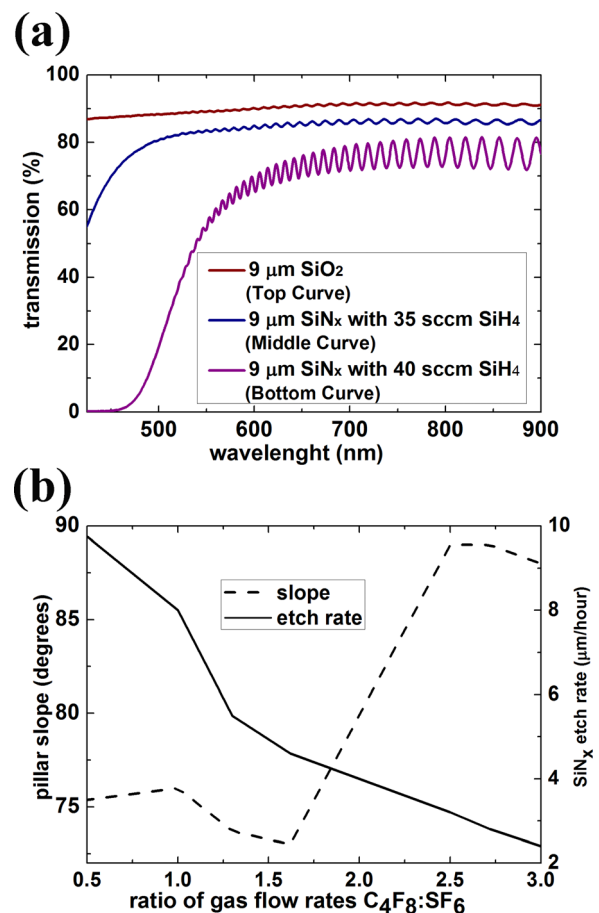


Fig. 3. (Color online) (a) Measured optical transmission spectra of the thin films used in light pipe fabrication. (b) Solid curve: Dependence of  $\text{SiN}_x$  etch rate upon ratio between  $\text{C}_4\text{F}_8$  and  $\text{SF}_6$  gases. Dotted curve: dependence of slope of light pipe walls with etch gas ratio.

power: 12 W, pressure: 10 mTorr, and temperature: 25 °C). The pillar is 8.5  $\mu\text{m}$  tall and has a diameter of 4  $\mu\text{m}$ . The diameters of the disks of the photomask are 5  $\mu\text{m}$ . Polymer produced during the etch process can be seen near the bottom of the lower light pipe of Fig. 4(a). In order to remove the undesirable polymer, we add O<sub>2</sub> gas (flow rate 4 sccm) to the SF<sub>6</sub>/C<sub>4</sub>F<sub>8</sub> ICP RIE process, and increase the pressure to 15 mTorr. Other parameters remain the same (C<sub>4</sub>F<sub>8</sub> flow

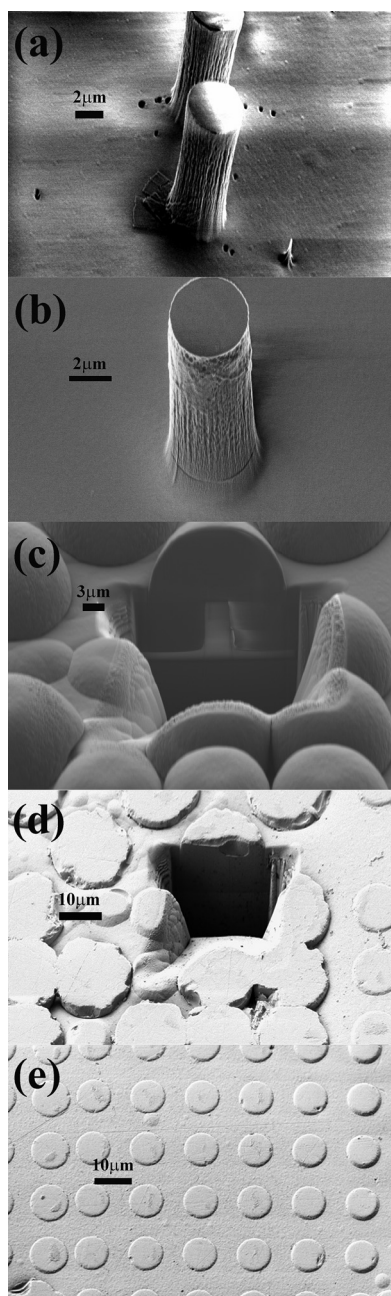


FIG. 4. (a) SEM image of SiN<sub>x</sub> light pipes with heights of 8.5  $\mu\text{m}$ . Dry etching is done in STS RIE system using an optimized recipe for vertical sidewalls and smooth side wall surfaces. Residual polymer can be seen at bottom of pipe. (b) SEM image of the SiN<sub>x</sub> light pipes after removal of Al etch mask disks. Various methods for reducing the formation of polymer are applied (see text). (c) SEM image of sample upon which FIB milling has been employed to produce cross section showing light pipe after thick SiO<sub>2</sub> deposition. (d) SEM image after CMP process. Light pipe can also be seen in the cross sectional view (e) SEM image after CMP process is applied to a test sample having light pipes with smaller diameters (2  $\mu\text{m}$ ).

rate: 150 sccm, SF<sub>6</sub> flow rate: 75 sccm, coil power: 1200 W, and platen power: 12 W). The SiN<sub>x</sub> deposited with an SiH<sub>4</sub> flow rate of 35 sccm is found to be smoother than that deposited with a SiH<sub>4</sub> flow rate of 40 sccm, possibly due to the lower deposition rate (11.6 nm/min compared to 18.8 nm/min). After the etching process, we immerse the substrate into boiling acetone with ultrasound agitation for 2 min, followed by cleaning in an O<sub>2</sub> plasma for 30 min. The plasma employs an O<sub>2</sub> flow rate of 45 sccm, coil and platen powers of 800 W and 50, respectively, and a chamber pressure of 35 mTorr. These steps result in vertical SiN<sub>x</sub> light pipes with smoother and cleaner surfaces, as evidenced in Fig. 4(b). The light pipes have height/diameter aspect ratios of 2.13 and 8.2 for diameters of 4  $\mu\text{m}$  and 0.9  $\mu\text{m}$ , respectively. We then deposit SiO<sub>2</sub> by PECVD, whose highly conformal nature results in dome structures being formed over the light pipes, with radii approximately equal to the SiO<sub>2</sub> thickness (10  $\mu\text{m}$ ). This thickness (of 10  $\mu\text{m}$ ) refers to positions away from the light pipes. To obtain a cross sectional view that shows the light pipes underneath the domes, we use a focused ion beam (FIB) on a test sample. The results are shown in the SEM image of Fig. 4(c). We then use CMP to planarize the domes. In our CMP process, we glue the sample containing the SiN<sub>x</sub> light pipes and SiO<sub>2</sub> domes to a metal holder using wax melted by heating to 160 °C. The metal holder weighs  $\sim$ 100 g. We then place the sample (with attached metal holder) face-down on a piece of lapping paper in a container filled with deionized (DI) water. The DI water acts as a lubricant. The lapping paper (3M Superabrasives and Microfinishing Systems Division; www.solutions.3m.com) comprises silicon carbide or aluminum oxide cylinders with diameters  $\sim$ 1–3  $\mu\text{m}$ . The metal holder (to which the sample is attached) is moved over the lapping paper by hand in a sequence of eight different patterns in order to polish the sample surface uniformly. This process is repeated for 30 min, after which a decrease in dome height of  $\sim$ 8  $\mu\text{m}$  is measured. The results are shown in the SEM of Fig. 4(d), in which an FIB is again used to produce a cross sectional view. For smaller light pipes (2  $\mu\text{m}$  diameter), the polishing process can be seen to produce a final result that is more planar [Fig. 4(e)]. This is because for these pipes, the distance between the domes is larger and there is less SiO<sub>2</sub> to remove. In both cases [Figs. 4(d) and 4(e)], it can be seen that the domes have been largely polished away. Portions of the domes remain, but the surface is sufficiently flat so that it is possible to test the devices in a way that demonstrates the improved optical performance enabled by the light pipes. The results demonstrate the successful fabrication of silicon nitride light pipes embedded in silicon dioxide. In the next section, we demonstrate their integration with silicon photodetectors.

## B. Photodetector fabrication and integration with silicon nitride light pipes

The photodiodes we report are lateral p-i-n structures. The starting substrate is an SOI wafer, with the [100] crystal orientation. The top layer is a 4  $\mu\text{m}$  thick lightly p-type

doped silicon film with  $1000 \Omega\text{-cm}$  resistivity, with a  $1 \mu\text{m}$  thick  $\text{SiO}_2$  insulator layer on top of a  $500 \mu\text{m}$  thick Si wafer. We begin by depositing silicon nitride to a thickness of  $40 \text{ nm}$  on the wafer to prevent oxidation and contamination. Alignment markers are defined by photolithography. Shipley 1813 photoresist with a thickness of  $\sim 1.3 \mu\text{m}$  is used. A resolution of  $\sim 1 \mu\text{m}$  is obtained by vacuum contact mode lithography (KarlSuss MJB3 or MJB4 mask aligners). The alignment markers are then etched into the substrate to a depth of  $\sim 4 \mu\text{m}$  by reactive ion etching (RIE) with  $\text{CF}_4$  and Ar gases. We next form mesas, to isolate the photodiodes from one another electrically. This involves photolithography, this time using SU-8 resist (2002,  $\sim 2 \mu\text{m}$  thick) as the etch mask. We then etch to a depth  $4.2 \mu\text{m}$ . The top silicon is therefore completely etched through to the  $\text{SiO}_2$  insulating layer, ensuring electrical isolation of the photodiodes. We employ a release layer (Omnicore from Micro Chem) with the SU-8 resist. This facilitates removal of the SU-8, which can otherwise be hard to achieve. The Omnicore release layer and SU-8 are removed with an  $\text{O}_2$  plasma. The wafer is then dipped into hydrofluoric acid (HF) to remove the oxide formed during the oxygen plasma process. We next produce the n+ type doped region of the photodiode. We first deposit a layer of silicon nitride to a thickness of  $100 \text{ nm}$  by PECVD. Lithography is then performed to define the region to be doped n+. Shipley S1813 resist is used ( $\sim 1.3 \mu\text{m}$  thick) and baked after development at a temperature of  $115^\circ\text{C}$  for  $1 \text{ min}$ . We then etch the silicon nitride film by reactive ion etching (Nexx RIE); the recipe is as follows:  $\text{CF}_4$  flow rate is  $15 \text{ sccm}$ , Ar flow rate is  $6 \text{ sccm}$ , pressure is  $10 \text{ mTorr}$ , RF power is  $100 \text{ W}$ , and the microwave power is  $306 \text{ W}$ . The resultant  $\text{SiN}_x$  etch rate is  $100 \text{ nm/min}$ . In this step, the underlying silicon is also etched, to a depth of  $\sim 0.1 \mu\text{m}$ . We next dope the regions of the wafer in which openings in the silicon nitride have been formed. Spin-on-doping is used. A layer of phosphosilica dopant is spun on to a silicon wafer, on which a silicon nitride film had been deposited. The wafer is baked at  $200^\circ\text{C}$  for  $15 \text{ min}$ . We then place the samples to be doped on a plain silicon wafer. Pieces of silicon wafers are then added to this plain silicon wafer to act as spacers. We then place the wafer containing the doping material on top. This structure is then placed in a furnace (Lindberg) at  $950^\circ\text{C}$  for  $15 \text{ min}$ . The furnace has nitrogen flow of  $140 \text{ sccm}$  and a forming gas flow of  $300 \text{ sccm}$ . The structure is then removed from the furnace. The top wafer, containing the doping material, is then removed, and the structure is returned to the furnace for annealing for a further  $15 \text{ min}$ . The samples are then placed in diluted HF (1:24 HF/water) to remove the phosphosilicate glass formed on the silicon surface. A similar fabrication method is used to dope the p+ regions, with borosilica instead of phosphosilica. A variety of photodiode structures are defined, with active region diameters ranging from  $2 \mu\text{m}$  to  $5 \mu\text{m}$ . The next step involves establishing the electrical connections to the photodiodes. We perform photolithography (Shipley S1813 resist), then thermally evaporate chrome and gold to thicknesses of  $100 \text{ \AA}$  and  $500 \text{ \AA}$ , respectively, with a deposition

rate of  $0.5 \text{ \AA s}^{-1}$ . The lift-off in acetone with ultrasonic agitation is performed for  $1 \text{ min}$ . The samples are then cleaned in acetone, methanol, and isopropanol, followed by a dehydration bake at  $120^\circ\text{C}$  for  $2 \text{ min}$ . The samples are then annealed in a furnace (Lindberg) at  $450^\circ\text{C}$  for  $40 \text{ s}$  to ensure that ohmic contacts to the Si are produced. This completes the fabrication of the silicon photodetectors. A microscope image of a typical device at this step is shown in Fig. 5(a). The active region width is  $20 \mu\text{m}$ . The wide metal pads that connect to the n+ and p+ regions are  $200 \mu\text{m}$  wide.

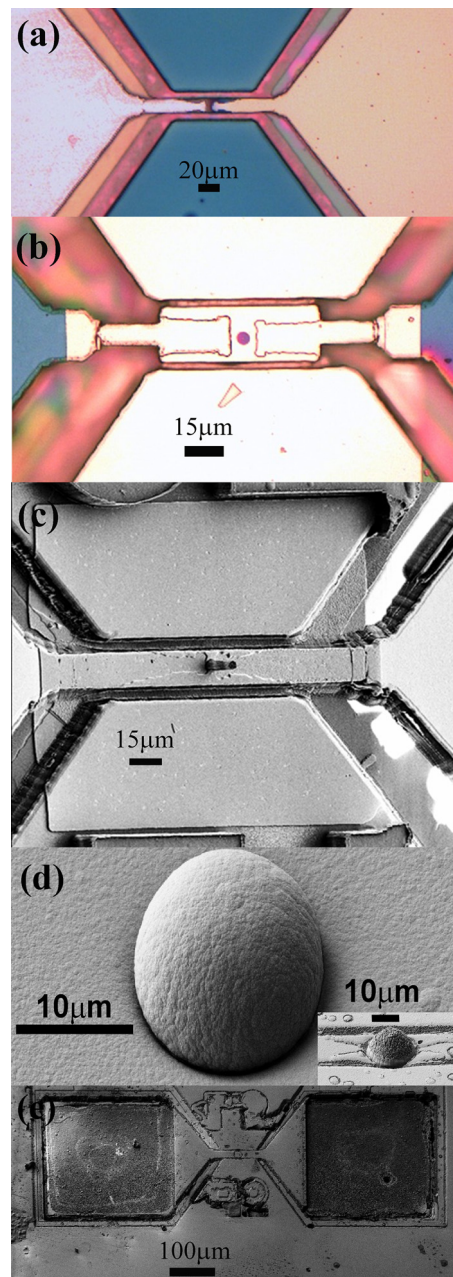


FIG. 5. (Color online) Photodiode fabrication. (a) Microscope image of complete lateral p-i-n Si-based photodiode (b) Microscope image of device after metal light blocking layer formation. (c) SEM image of  $\text{SiN}_x$  light pipe integrated with photodiode (d) SEM image of a test device which has only pipe after thick  $\text{SiO}_2$  deposition and inset shows the real device after thick  $\text{SiO}_2$  is deposited (e) SEM image after CMP process.

We now describe the fabrication of silicon nitride light pipes on the photodiodes. The process begins with the deposition of silicon nitride on the photodetectors to a thickness of 100 nm. This layer serves to electrically insulate the photodetectors from the metal layer we deposit next for light-blocking. It is therefore crucial that there be no pinholes in the silicon nitride film. The metal light-blocking layer serves to prevent light from being incident on the photodetector regions around the pillar, ensuring that the device has a photosensitive region with defined extent. The light-blocker layer is formed as follows. After the initial photolithography step, layers of chromium, aluminum, and chromium are evaporated to thicknesses of 20, 50, and 10 nm, respectively, followed by a lift-off process. The first Cr layer is for adhesion, while the second Cr layer serves to protect the Al when we later use a wet etchant (for Al) to remove the Al etch mask disks from the  $\text{SiN}_x$  light pipes. The resultant light-blocking layer appears as the bright region of Fig. 5(b). The opening in the light blocking layer, which represents the photosensitive region of the device, appears as a purple circle in Fig. 5(b). We next form the light pipes. Silicon nitride is deposited to a thickness of  $\sim 8 \mu\text{m}$  step using the process described in Sec. II A. The deposition parameters that yielded the silicon nitride with optimized optical properties (Fig. 3 with  $\text{SiH}_4$  flow rate of 35 sccm) are employed. We then define Al disks on the silicon nitride layer using photolithography, evaporation and lift-off. The disks are aligned with the active regions of the photodiodes. Light pipes are then produced by dry etching (ICP RIE), using the recipe described in Sec. II A. The light blocking metal layer also acts as an etch stop. The results are shown as Fig. 5(c). We next deposit  $\text{SiO}_2$  (borophosphosilicate glass) by PECVD to a thickness of  $\sim 9 \mu\text{m}$ . The results, obtained on a test device, are shown as Fig. 5(d). Photolithography is then performed to define the openings to be formed in the  $\text{SiO}_2$  that enable electrical connections to the photodetectors. Here, an image reversal photoresist (AZ5214E photoresist, Clariant) is used in order to achieve a negative wall profile, as this is advantageous for the lift-off process. Thermal evaporation of aluminum, chromium, and titanium layers is then performed to thicknesses of 300, 100, and 100 nm, respectively. The lift-off process is then carried out. In the following RIE step, the flow rates of the  $\text{C}_4\text{F}_8$  and  $\text{SF}_6$  gases are both 10 sccm, the pressure is 7 mTorr, and coil and platen powers of 800 W and 100 W are used. The etching is performed for 40 min, resulting in an etch depth of  $9.2 \mu\text{m}$ , as determined from a test sample. It can be seen from Fig. 5(d) inset that the  $\text{SiO}_2$  surface is relatively rough. The metal mask thickness is not uniform and becomes completely etched through in some regions, including the  $\text{SiO}_2$  domes, which then become rough due to micromasking. This is not a major source of concern, however, as the next step is planarization by CMP. This is motivated mainly to planarize the  $\text{SiO}_2$  domes that form over each silicon nitride light pipe, but smoothes the other surfaces as well. The result is shown as Fig. 5(e). We then perform wire bonding, which completes the device. A schematic depiction of the final device is included as Fig. 6.

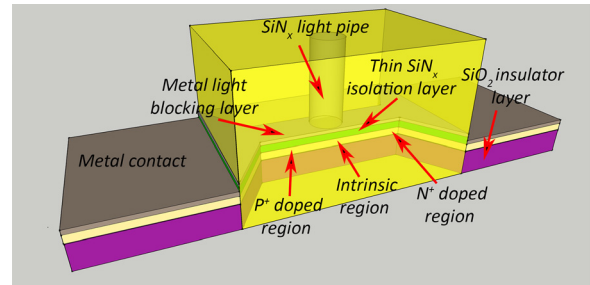


Fig. 6. (Color online) Schematic representation of completed device.

### III. MEASUREMENTS

#### A. Current–voltage measurements

Upon the completion of the microfabrication of the lateral Si-based p-i-n photodiodes, we perform current–voltage (I-V) measurements upon them under dark and illuminated conditions. Illumination is provided with light from the tungsten lamp of a probe station microscope. Measurements are performed with a Keithley 2400 electrometer and Signatone microprobes. The measurement data are shown in Fig. 7. It can be seen that, for bias voltages close to 0 V, the ratio of photocurrent (with microscope lamp illumination) to dark current is four orders of magnitude ( $3 \times 10^4$ ). The dark current is on the order of a few picoamperes, while the photocurrent is on the order of tens of nanoamperes.

#### B. Scanning photocurrent measurements

We next describe the scanning photocurrent measurements employed to confirm the improved optical performance enabled by the light pipe structures.<sup>5</sup> Measurements are made both on devices containing light pipes, and on light pipe-free devices. Microscope images of the devices are shown in Figs. 8(a) and 8(b). We first place the devices on a probe station equipped with an electrometer (Keithley 2400) and measure the dark current to be  $\sim 0.3 \mu\text{A}$  at a bias voltage of  $-0.2 \text{ V}$  for both types of devices. We use this bias voltage in the scanning photocurrent measurements we perform

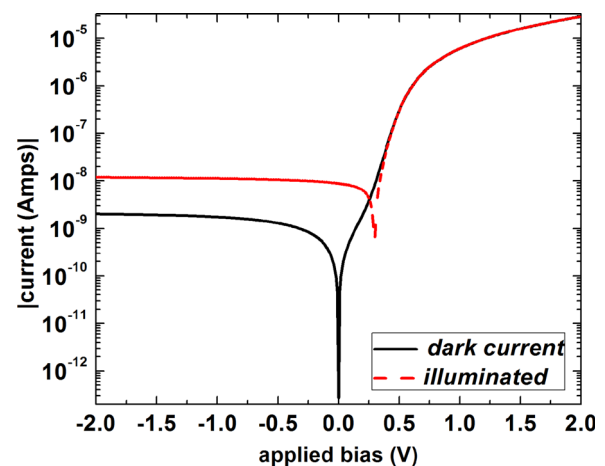


Fig. 7. (Color online) I-V characteristics of the fabricated Si based lateral p-i-n type photodetector. Photocurrent to dark current ratio at 0 V bias voltage is 4 orders of magnitude.

next. Each device is placed in a sample-scanning confocal microscope (WiTEC). In this system, fiber-coupled laser light ( $\lambda = 532$  nm) is collimated, passed through a chopper, and focused by an objective lens (NA = 0.9, magnification = 100 $\times$ ) onto the device, which sits on a piezoelectric translation stage. The laser power from the objective lens is  $\sim 7$   $\mu$ W. The device is biased with an electrometer (Keithley 2400), and the current is measured with a lock-in amplifier. The reference signal is provided by the chopper, enabling the photocurrent to be extracted from the total current. The photocurrent is then recorded as a function of position as the device is scanned. Results obtained in this way are shown as Figs. 8(c) and 8(d) for the light pipe device and light pipe-free device, respectively. The bright circular spot of each image corresponds to the case where the focused laser beam is centered over the light pipe entrance. It should be noted that this data is obtained for the case where the laser spot is focused on the top surface. We observe that the photocurrent increases when the laser spot is centered over the mesa isolation trenches, but this does not reflect on the analysis of the light pipe properties. From Figs. 8(c) and 8(d), it can be seen that considerably higher photocurrent results from the light

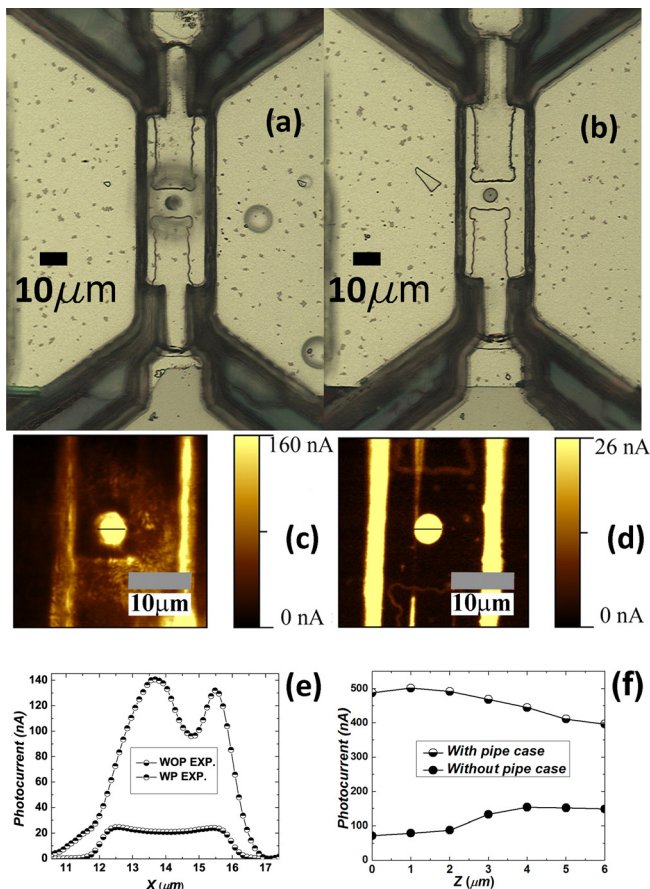


FIG. 8. (Color online) (a) Optical microscope image of completed device with light pipe. (b) Optical microscope image of light pipe free device fabricated for comparison purposes. Scanning photocurrent microscopy maps of (c) light pipe and (d) light pipe-free devices. (e) Photocurrent vs  $x$ -distance for light pipe (“experiment with light pipe”) and light pipe-free (“experiment with out light pipe”) devices measured along paths indicated in panels (c) and (d) with black lines. (f) Photocurrent vs vertical sample position ( $z$ ), for another pair of light pipe and light pipe-free devices, measured under higher laser power.

pipe device than the light pipe-free device. To explore this further, in Figs. 8(e) and 8(f) we plot the photocurrent as a function of distance for the devices with and without light pipes along the cross sections indicated in Figs. 8(c) and 8(d) (black lines). Figures 8(e) and 8(f) show that the peak photocurrent for the light pipe device is  $\sim 6.3$  times larger than the peak value of the light pipe-free device. Interestingly, local minima appear in the centers of the photocurrent profiles for both the device containing the light pipe, and for the light-pipe free device. The cause of these minima is not fully understood. If we compare the photocurrent within these local minima, the ratio becomes  $\sim 4.2$  times. We next consider the physical interpretation for the improvement enabled by the light pipe. Rather than full-field numerical electromagnetic modeling, we employ a ray-tracing model. While this is not as rigorous, it provides helpful physical insight. We note that our method is appropriate for the waveguides we study, as their diameters are significantly larger than the laser wavelength. For smaller waveguides, full-field simulations (e.g., Ref. 2) should be used. Full details of the method are provided in supplementary material document of Ref. 5, and only summarized here. We consider the light focused on the top surface of  $\text{SiO}_2$  as consisting of a collection of rays with angles ranging from 0 to  $\theta_{\text{NA}}$ , where  $\text{NA} = \sin \theta_{\text{NA}}$  is the numerical aperture ( $=0.9$ ) of the microscope objective. Ray tracing is then employed to determine the largest angle  $\theta_{\text{max}}$  that a ray can be incident upon the top  $\text{SiO}_2$  surface, and still be collected by the silicon photodetector, whose diameter ( $R = 5$   $\mu\text{m}$ ) is defined by the opening in the Al light blocking layer. The light pipe device is regarded as comprising a 7  $\mu\text{m}$  tall  $\text{SiN}_x$  cylinder (diameter 5  $\mu\text{m}$ ) encased in  $\text{SiO}_2$  (8.7  $\mu\text{m}$  thick). The light pipe-free device contains just the 8.7  $\mu\text{m}$  thick  $\text{SiO}_2$ . The ray tracing predicts  $\theta_{\text{max}} \approx 64^\circ$  for the light-pipe device, while  $\theta_{\text{max}} \approx 24^\circ$  for the light-pipe free device. The photocurrents for the light pipe device, and the light pipe free device, are then predicted by integrating the intensities of the rays from 0 to  $\theta_{\text{max}}$ , but with the reflection coefficients at the interface (air- $\text{SiO}_2$ ,  $\text{SiO}_2$ - $\text{SiN}_x$ ,  $\text{SiO}_2$ -Si, and  $\text{SiN}_x$ -Si) taken into account. For simplicity, these coefficients are found by averaging the Fresnel reflection coefficients for p- and s-polarizations. The results predict that a photocurrent for the light pipe device that is  $\sim 3.6$  times larger than the light pipe-free device, in agreement with the experimentally measured trend.

To gain further insight into the improvement provided by the light pipes, we perform photocurrent measurements in which the device is scanned in the vertical ( $z$ ) direction. The results are obtained by centering the laser spot over the light pipe, or over the photodetector for the light pipe free device, then measuring the photocurrent as the device is vertically translated by the piezoelectric stage over a distance of 6  $\mu\text{m}$ . The results [Fig. 8(f)] show higher photocurrents than before due to the laser power being increased to 25  $\mu\text{W}$ . In this figure,  $z = 0$   $\mu\text{m}$  corresponds to the laser spot being focused at the  $\text{SiO}_2$  surface, while  $z = 6$   $\mu\text{m}$  corresponds to the device being moved by 6  $\mu\text{m}$  so that the laser spot is within the device. It can be seen that, for the light pipe device, the photocurrent maximum is reached at  $z \approx 1$   $\mu\text{m}$ , due to the



coupling to the light pipe being largest when the laser beam is focused at its entrance. It then decreases as  $z$  increases, with the photocurrent being  $\sim 1.3$  times higher at  $z = 0 \mu\text{m}$  than  $z = 6 \mu\text{m}$ . This is consistent with the value predicted by the ray-tracing analysis ( $\sim 1.2$  times). The photocurrent of the light pipe-free device increases  $\sim 1.6$  times as the position is varied from  $z = 0 \mu\text{m}$  to  $z = 6 \mu\text{m}$ , due to the waist being closer to the photodetector. This is again consistent with the trend predicted by ray-tracing of an increase of  $\sim 2.3$  times.

#### IV. CONCLUSION

In conclusion, we have demonstrated a fabrication method for vertical silicon nitride waveguides integrated with silicon photodetectors. We show that thick silicon nitride films with good optical transmission in the visible wavelength range can be achieved by optimization of the deposition conditions. We demonstrate that the silicon nitride pillars with near-vertical sidewalls can be achieved by fine tuning of the ratio between the flow rates of the etch gases. We also present that these silicon nitride pillars can be encased in  $\text{SiO}_2$  by PECVD, followed by a simple CMP process. The final integration of the light pipes with silicon photodetectors is also described, by utilizing a new technique that enables the photosensitive region of the photodetectors to be well defined in its spatial extent. This involves the deposition of a light blocking layer (Al-Cr), with an opening defined in its center. To characterize the finished

light pipe and light pipe-free devices, we perform current—voltage and scanning photocurrent measurements. The latter clearly demonstrate the improvements enabled by the light pipes, and are found to be consistent with ray-tracing modeling. We therefore expect that the incorporation of light pipes into CMOS image sensors would result in significantly increased efficiency and reduced cross-talk.

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